

A Theory Perspective on Axions

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at the BREAD Collaboration Meeting, Jan. 14, 2026

Outline

Key things to know about axions and Strong CP

- Demystifying theta terms
- Strong CP problem and axion solutions
- What is the right experimental target?

Axion theory at Harvard

- Brief summary: what my students and I are up to

The θ Parameter in $SU(n)$ Gauge Theory

Gauge theories have a θ **parameter**. There are formulas for it:

Lagrangian viewpoint:

$$S_\theta = S_0 + \frac{\theta}{8\pi^2} \int \text{tr}(F \wedge F) = S_0 + \frac{\theta}{64\pi^2} \int d^4x \epsilon^{\mu\nu\rho\sigma} F_{\mu\nu}^a F_{\rho\sigma}^a.$$

Hamiltonian viewpoint, Weyl gauge ($A_0 = 0$):

$$H_\theta = g^2 \left(i \frac{\delta}{\delta A_i^a} - \frac{\theta}{4\pi^2} B_i^a \right)^2 + \frac{1}{g^2} (B_i^a)^2 \quad \text{where} \quad B_i^a = (\nabla \times A^a)_i - \frac{1}{2} f^{abc} (A^b \times A^c)_i$$

Facts About the θ Parameter

Physical quantities depend on the value of θ .

All physics is 2π periodic in the value of θ . More precisely,

$$H_{\theta+2\pi} = UH_{\theta}U^{\dagger}$$

where U is a unitary operator (the “Chern-Simons operator”).

If θ is *not* an integer multiple of π , then **CP (equivalently, **T**) symmetry is explicitly violated.**

The θ term is “topological”: $\frac{1}{8\pi^2} \int \text{tr}(F \wedge F) \in \mathbb{Z}$ when integrating over a closed spacetime. (Mathematically, related to the 2π -periodicity.)

Facts About the θ Parameter

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Theorists like to emphasize this bit about the topological nature of the term, but people often get confused about it.

In particular, the other facts that I told you are **all true even if you study QFT in topologically trivial spacetime (like Minkowski spacetime)**.

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Aside: Advice About the θ Parameter

If you ever get confused about the θ parameter (many of us do, sometimes), there is a simple system with all the key properties.

This is a *free* quantum particle on a ring (or circle), with coordinate $\varphi \cong \varphi + 2\pi$, and action

$$S = \int dt \left(\frac{I}{2} \dot{\varphi}^2 + \frac{\theta}{2\pi} \dot{\varphi} \right)$$

It has Hamiltonian

$$H = \frac{1}{2I} \left(\frac{d}{d\varphi} - \frac{\theta}{2\pi} \right)^2$$

and θ -dependent energy eigenvalues $E_n = \frac{1}{2I} \left(n - \frac{\theta}{2\pi} \right)^2, \quad n \in \mathbb{Z}.$

Any surprising claim about θ terms can be evaluated in this context.

The Wrong Way to Approach the θ Parameter

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For *any* QFT with *any* parameter y , you can make up a kind of “super-QFT”^{*} with states $|\psi, y\rangle$ and matrix elements

$$\langle\psi', y'| O | \psi, y\rangle = \delta_{yy'} \left(\langle\psi'| O | \psi\rangle \right)_{\text{QFT with param } y}$$

^{*} The correct technical term is “direct sum of QFTs.”

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You *can* do this, but it’s usually a silly thing to do.

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When people talk about “ θ vacua,” they are making up such a “super-QFT.”

Oddly, people only do this for θ ; you never hear about, say, “electron Yukawa vacua,” but the super-QFT works there too.

People often say we do this because the “ θ vacua” are linear combinations of something else called “ n vacua,” but those make less sense.

What happened here is that a few very smart and influential people were confused in 1976, and other people have been repeating their confusion ever since.

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Fifty years have passed. We should stop repeating them.
 θ is *just* a parameter in the Lagrangian (or Hamiltonian).
Don't mystify it.

Neutron EDM and Strong CP Puzzle

Smaller than geometric cartoon expectation
by factor $m_u/\Lambda_{\text{QCD}} \sim 10^{-3}$. **(Chiral anomaly)**

The Strong CP phase $\bar{\theta}$ gives rise directly to a CP-odd pion-proton-neutron interaction.

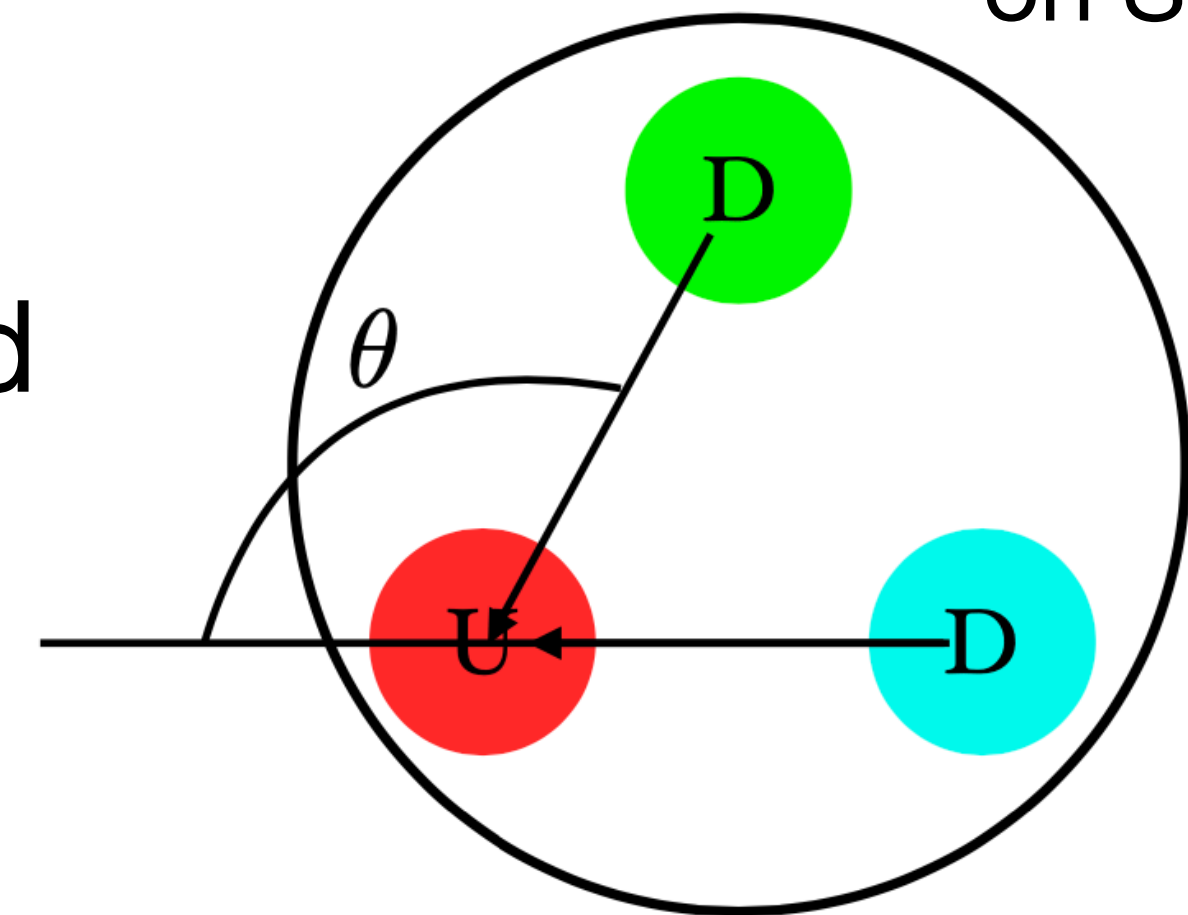
This in turn feeds into the neutron EDM at one loop:

$$d_n \approx \frac{e\bar{\theta}g_A c_+ m_u m_d}{8\pi^2(m_u + m_d)f_\pi} \log \frac{m_\rho^2}{m_\pi^2} \sim 3 \times 10^{-16} \bar{\theta} e \text{ cm} \Rightarrow |\bar{\theta}| \lesssim 10^{-10} (!)$$

Crewther, di Vecchia, Veneziano, Witten 1979

But: CP is not a symmetry of nature! Why so small?

Fig. from
Anson Hook,
TASI Lectures
on Strong CP



Axions

Promote θ to a **dynamical field**, $\theta(x)$, interacting with gluons.

$$\mathcal{L} = \frac{1}{2} f^2 (\partial\theta)^2 + \frac{1}{64\pi^2} \theta(x) G_{\mu\nu}^a \tilde{G}^{a\mu\nu}(x)$$

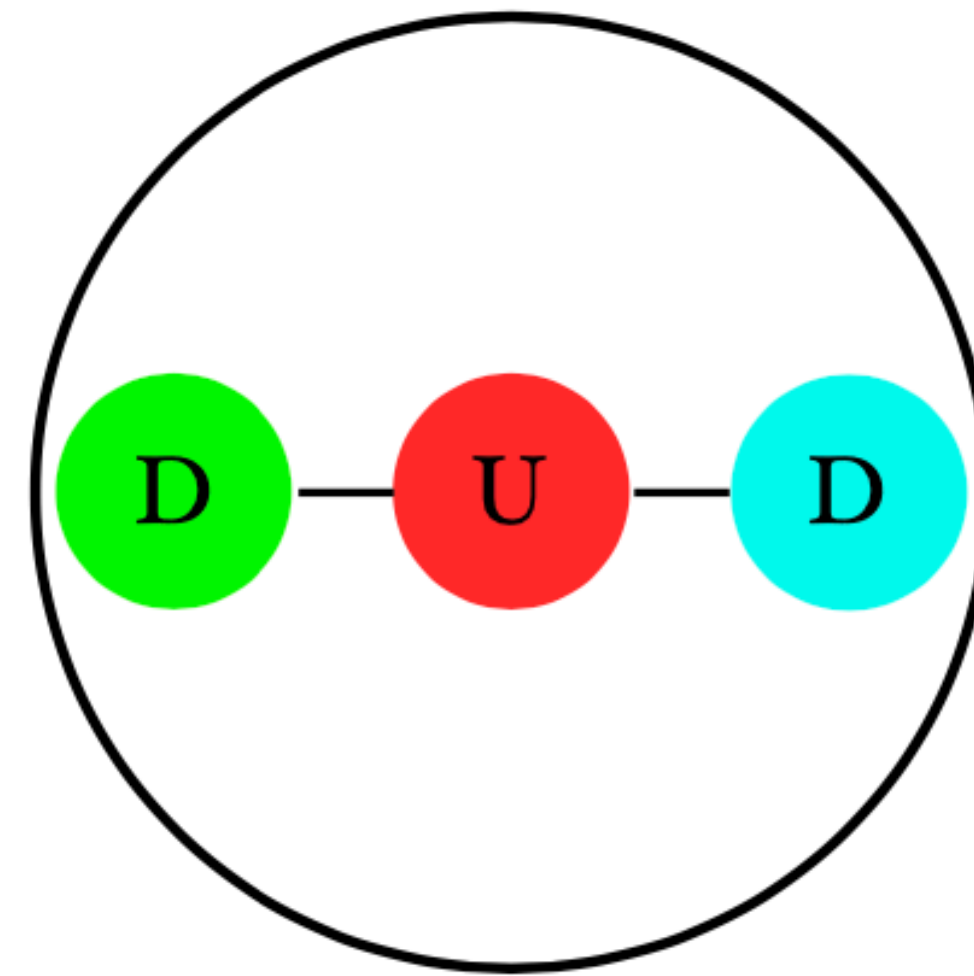
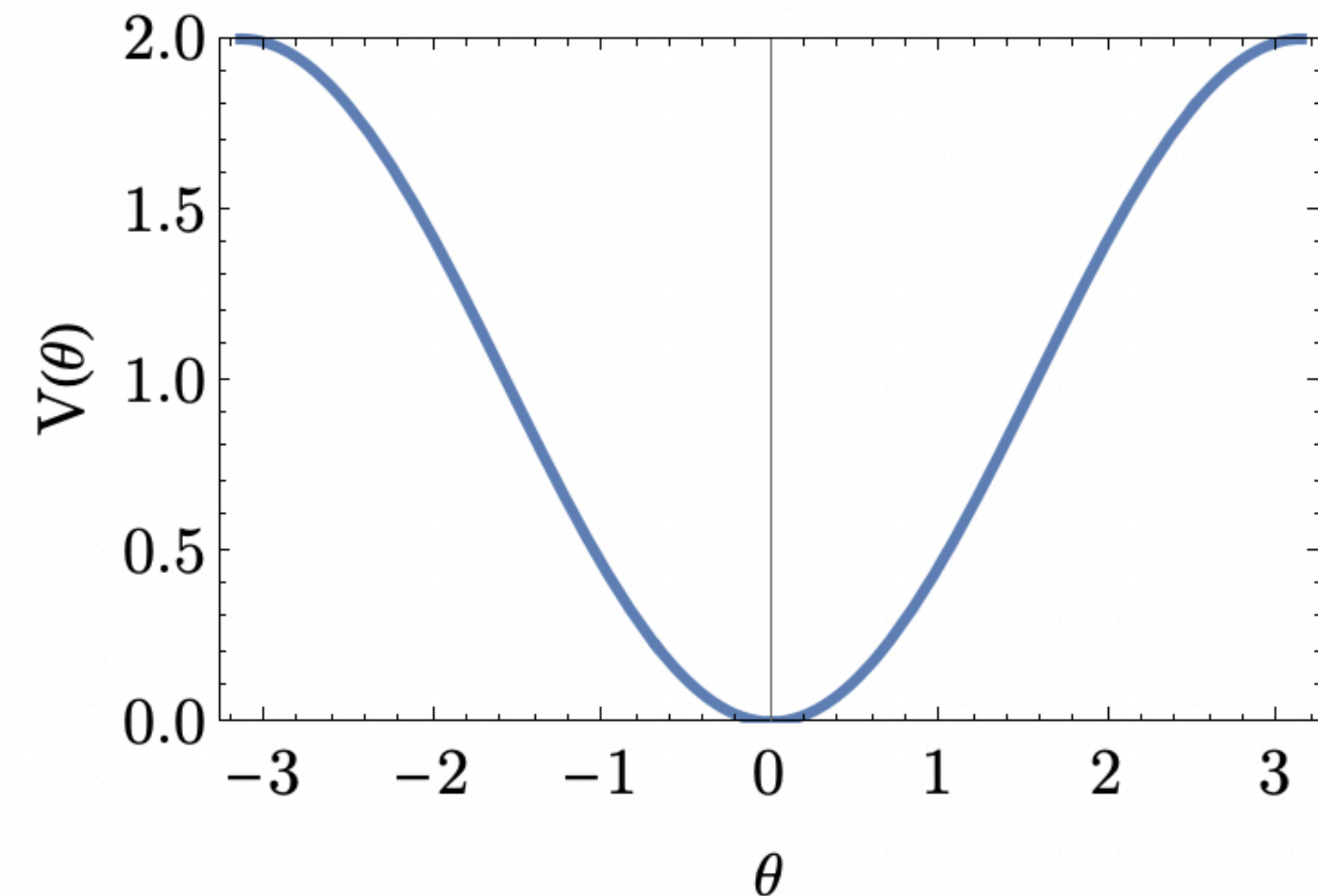


Fig. from Anson Hook, TASI Lectures on Strong CP



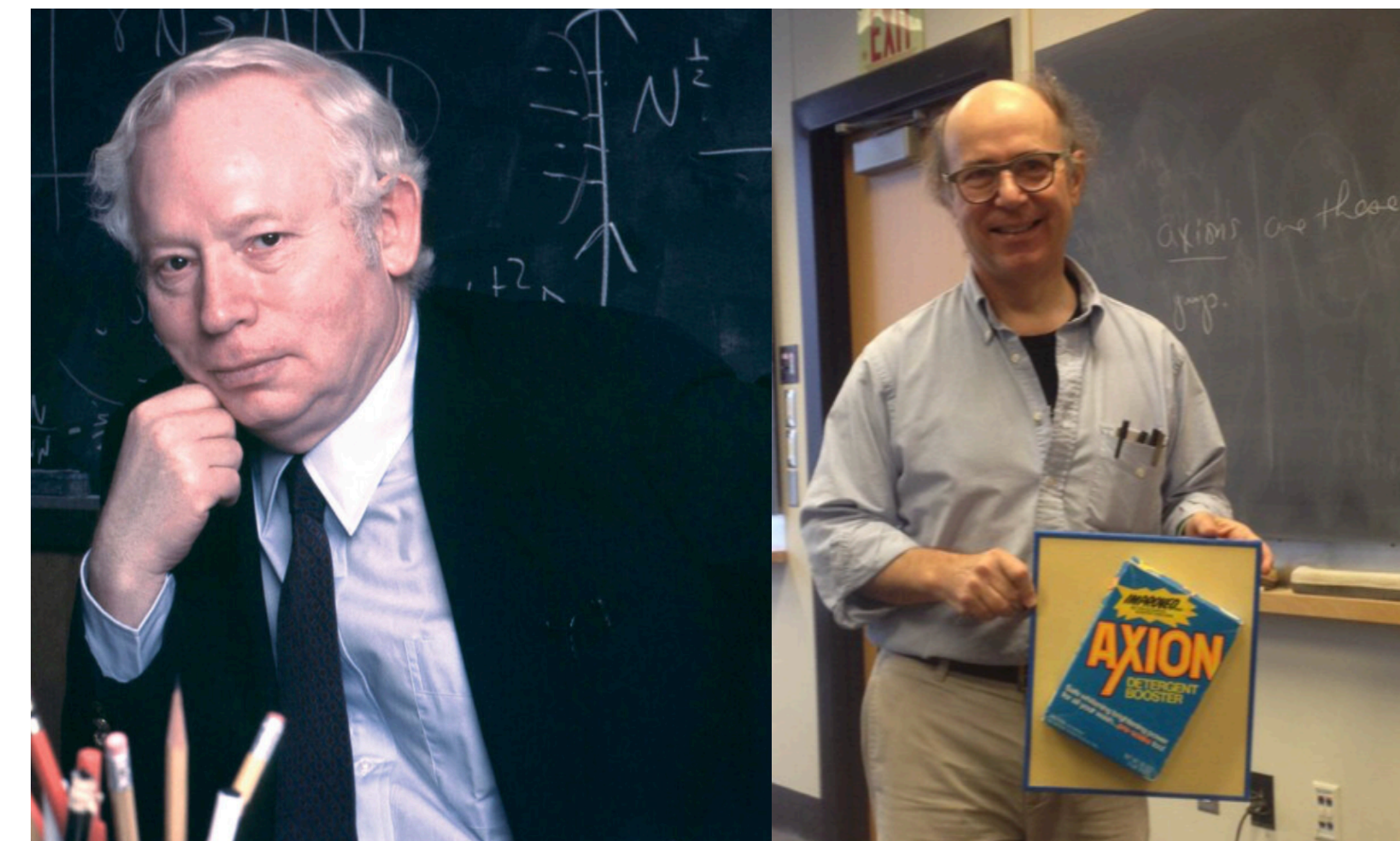
Roberto Peccei, Helen Quinn
(photo: Ryan Schude, Quanta Magazine)



Strong dynamics \Rightarrow

**Axion relaxes to
CP-conserving value.**

$$m_a \sim \frac{m_\pi f_\pi}{f} \sim \frac{10^{-5} \text{ eV}}{f/(10^{12} \text{ GeV})}$$

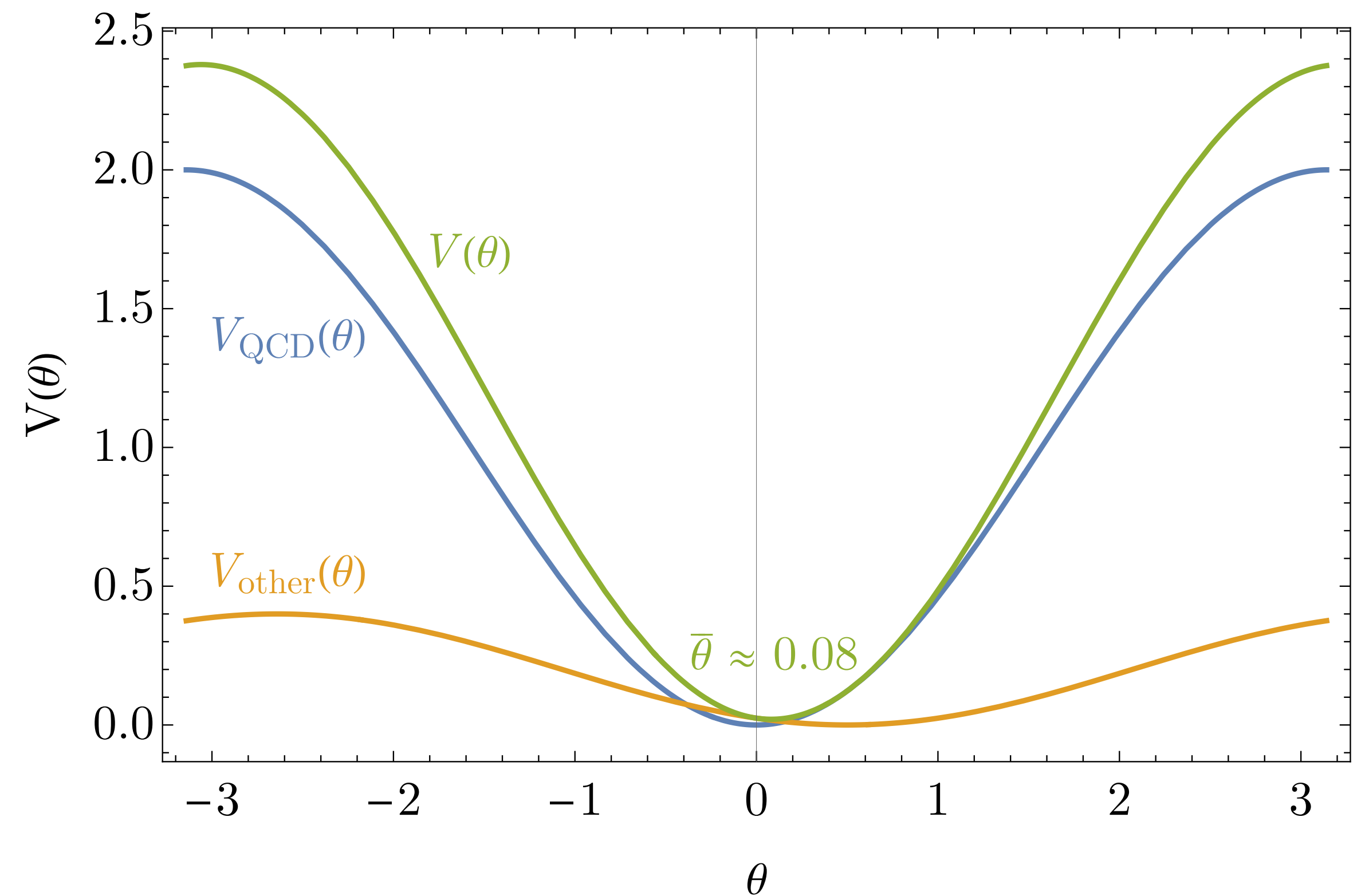


Steven Weinberg, Frank Wilczek

The Axion “Quality Problem”

The Peccei/Quinn (and KSVZ, DFSZ) “solution” wasn’t really. $V_{\text{QCD}}(\theta)$ is minimized at $\theta = 0$, but we could add all kinds of other terms that have a minimum elsewhere.

To explain $|\bar{\theta}| < 10^{-10}$ we need to assume $V_{\text{other}}(\theta)$ to be tiny — but then why not just assume small $\bar{\theta}$ to start?



(Georgi, Hall, Wise '81; Lazarides, Panagiotakopoulos, Shafi '86; Casas, Ross '87; Kamionkowski, March-Russell '92; Holman, Hsu, Kephart, Kolb, Watkins, Widrow '92; Barr, Seckel '92; Ghigna, Lusignoli, Roncadelli '92; ...)

Axion Theories: 4d (“Peccei-Quinn”) vs. Extra-Dimensional

Familiar 4d axion theories (e.g., KSVZ-like, DFSZ-like): the axion is a **pseudo-Nambu-Goldstone boson**, the phase of a complex field Φ . Quality problem from terms like $\frac{\lambda}{M_{\text{Pl}}^{n-4}} \Phi^n + \text{c.c.}$ in the Lagrangian, which must be suppressed all the way up to $n \sim 14$. **New gauge symmetries** can do that.

Extra-dimensional axion theories: the axion is **not** a Goldstone boson.

Arises as gauge field **integrated over extra dimension(s)**: $\theta = \int A$. Gauge invariance forbids terms like A^n , allows terms involving $F = dA$ hence $d\theta$.

Experimentally: hard to tell the difference. (Maybe best hope is cosmology.)

Some Good News

For **extra-dimensional axions**, the quality problem is mild. Every dangerous term $V_{\text{other}}(\theta)$ is proportional to an exponential like $\exp(-M^d L^d)$ where M is a fundamental mass scale and L a length scale of extra dimensions.

Axion-gluon coupling from a “Chern-Simons” coupling $\frac{1}{8\pi^2} \int A \wedge \text{tr}(G \wedge G)$.

Not *guaranteed* to solve the Strong CP problem, but often will.

Furthermore, **extra-dimensional axions** are a **prediction** of string theory. They are probably the single most promising way we might connect quantum gravity with realistic experiments in the foreseeable future.

(This all goes back to a 1984 paper of Edward Witten; see also 2006 work of Joe Conlon and of Peter Svrcek & Edward Witten.)

Axion Models at a Glance

Pseudo-Nambu-Goldstone
for 4d $U(1)_{PQ}$

Zero mode of gauge field in
higher dimensions

“Pre-inflation”
scenario

Quality problem

(Quality problem)

Isocurvature problem

Isocurvature problem

UV insights

Post-inflation
PQ transition

Quality problem

Domain wall problem

Stable relic problem

Not possible
(no linearly realized PQ
symmetry to break)

(maybe similar late-time physics from
other initial conditions?)

Axion interactions with gluons & photons

Fundamental axion interactions:

$$\int \left(\underbrace{\frac{k_G}{8\pi^2} \theta(x) \operatorname{tr}(G \wedge G)}_{\text{(gluons)}} + \underbrace{\frac{k_F}{8\pi^2} \theta(x) F \wedge F}_{\text{(photons)}} \right)$$

The interactions are **quantized**. If no fractionally charged particles:

$$k_G \in \mathbb{Z}, \quad \frac{2}{3}k_G + k_F \in \mathbb{Z}$$

(MR '23; Choi, Forslund, Lam, Shao '23; Córdova, Hong, Wang '23)

Low-energy axion-photon coupling: $g_{a\gamma\gamma} = \frac{\alpha}{2\pi f/k_G} \left(\frac{2k_F}{k_G} - 1.92(4) \right)$

Axion interactions with gluons & photons

Constrained by
quantization

$$k_G \in \mathbb{Z}, \quad \frac{2}{3}k_G + k_F \in \mathbb{Z}$$

Low-energy axion-
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$$g_{a\gamma\gamma} = \frac{\alpha}{2\pi f/k_G} \left(\frac{2k_F}{k_G} - 1.92(4) \right)$$

From mixing
with pions

A minimal $|k_G| = 1$ is often theoretically favored (e.g., domain wall problem).

Then the **smallest** $g_{a\gamma\gamma}$ is when

$$\frac{2k_F}{k_G} \equiv \frac{E}{N} = \frac{8}{3}.$$

a.k.a.
“DfSZ target”
but **not** linked
to DfSZ model

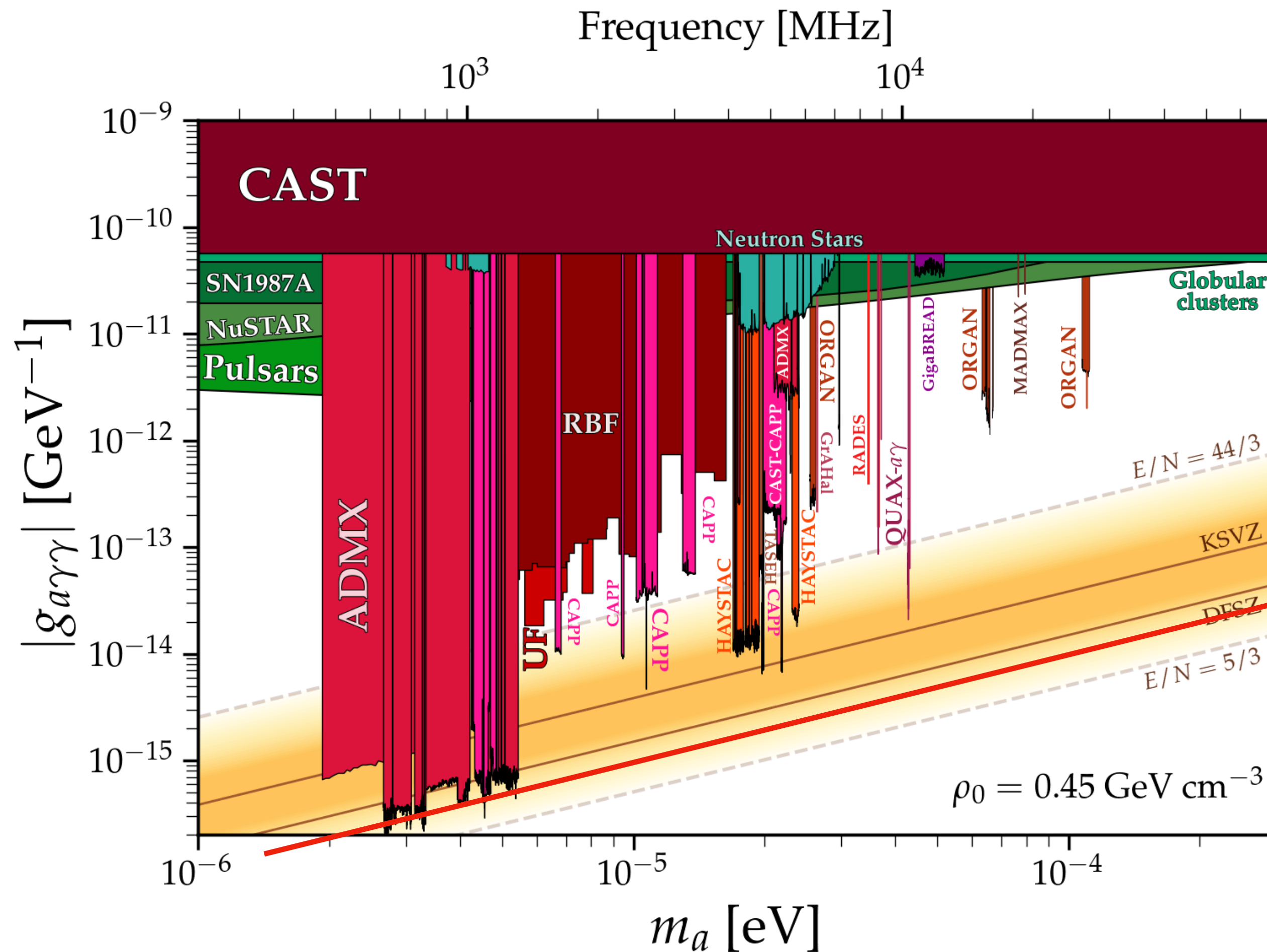
Experimental Target

Experiments measure $\rho_0 |g_{a\gamma\gamma}|^2$.

“DFSZ line” $E/N = 8/3$ **with**
 $\rho_0 = 0.45 \text{ GeV/cm}^3$ **dark**
matter density.

Credible estimates have large
 error bars (e.g., de Salas &
 Widmark '20), often centered
 at **smaller** ρ_0 , so need to reach
lower in the plot.

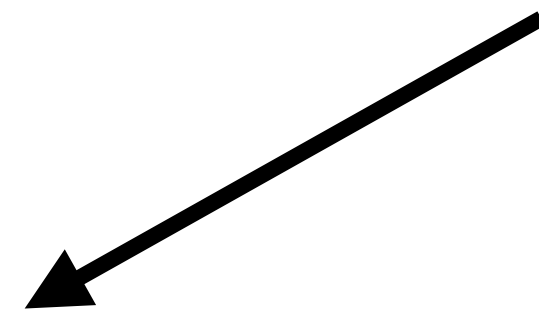
My nightmare scenario is we
 stop looking because “it’s
 already ruled out” at **wrong** ρ_0 .



Source: Ciaran O’Hare, github

Some Ongoing Research Directions

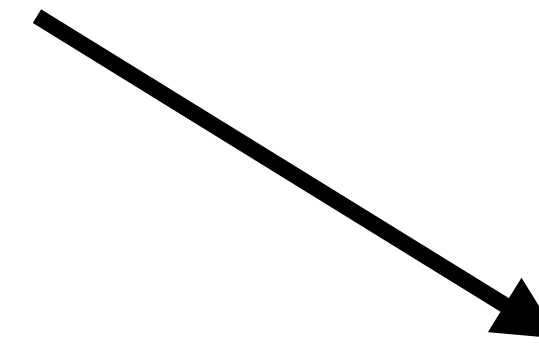
Reducing axion isocurvature for an extra-dimensional, pre-inflation axion:
axion properties different today than during inflation.



Time-dependent mass, first-order phase transition, gravitational waves,



w/ Prish Chakraborty, Junyi Cheng, Zekai Wang
arXiv:2507.12519 and ongoing work



Time-dependent decay constant, moduli dynamics, reheating, ...



w/ Chandrika Chandrashekar (ongoing work)

And: chiral fermions \Rightarrow existence of light axions in string theory

Brief Advertisement

Local students in the audience, please consider taking my class:

Axion Physics

PHYSICS 253DR Matthew Reece ▼

Class Number: 17222 Course ID: 226491 Consent: No Consent Class Capacity: No Limit

View:

Description

Q Scores/Course Evaluations

Theory and phenomenology of axion fields. First, we review the QFT toolkit needed to understand the Strong CP problem and axions (e.g., instantons, theta terms, Chern-Simons terms, the chiral anomaly, the chiral Lagrangian). We will discuss axion models with an emphasis on the axion quality problem and extra-dimensional axions; phenomenological aspects of axion physics (axion cosmology, axion dark matter, laboratory searches, astrophysical constraints); and theoretical developments related to axions (anomaly inflow, generalized symmetries).

Recommended Prep: Familiarity with quantum field theory and the Standard Model (e.g., Physics 253a, 253b and/or 254) is necessary. Some familiarity with cosmology would be useful, but we will review the basics.

2026 Spring
Full Term
1/26/2026 to 4/29/2026
S M T W Th F S
1:30pm - 2:45pm
Location:TBA3

Thank You!